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# The pressure-induced phase transition of Ga<sub>2</sub>O<sub>3</sub>

Baozhao Tu<sup>1</sup>, Qiliang Cui<sup>1,3</sup>, Peng Xu<sup>1</sup>, Xin Wang<sup>1</sup>, Wei Gao<sup>1</sup>, Chengxin Wang<sup>1</sup>, Jing Liu<sup>1,2</sup> and Guangtian Zou<sup>1</sup>

 <sup>1</sup> National Laboratary of Superhard Materials, Jilin University, Changchun 130023, People's Republic of China
<sup>2</sup> BSRF, Institute of High Energy Physics, Academia Sinica, Beijing 100039, People's Republic of China

E-mail: qiliangcui@yahoo.com.cn

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#### Abstract

In this paper, we investigate the pressure-induced structural transition of  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub> powder by means of energy-dispersive x-ray diffraction (EDXD) measurements. The EDXD results show that, with increase of pressure, a new pressure-induced phase appears. The new tetragonal structure ( $\beta$ -Ga<sub>2</sub>O<sub>3</sub>) can remain stable over the pressure range ( $\leq$ 38 GPa) under study.

### 1. Introduction

 $Ga_2O_3$  is a wide-band-gap ( $E_g = 4.9 \text{ eV}$ ) [1] compound which has long been known to exhibit both conduction and luminescence properties [2–4]. Nanoscale  $Ga_2O_3$ , which is made from  $Ga_2N$ , also been researched in depth. But its properties under high pressure have never been reported. In this paper, we report the properties of the pressure-induced phase transition of  $Ga_2O_3$ .

Ga<sub>2</sub>O<sub>3</sub> can exists in four types; they are  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub>,  $\beta$ -Ga<sub>2</sub>O<sub>3</sub>,  $\gamma$ -Ga<sub>2</sub>O<sub>3</sub> and  $\varepsilon$ -Ga<sub>2</sub>O<sub>3</sub>. The most stable one is  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub>, which is monoclinic.  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> and  $\gamma$ -Ga<sub>2</sub>O<sub>3</sub> are harder materials. The  $\beta$ -type material has the tetragonal structure and the  $\gamma$ -type form shows the behaviour of a cubic structure.

#### 2. Experiment

Our sample is white powder of Ga<sub>2</sub>O<sub>3</sub>. The XRD results under ambient pressure and temperature show that the sample is an  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub> crystal. Figure 1 shows the energy-dispersive x-ray diffraction spectra of Ga<sub>2</sub>O<sub>3</sub> at ambient pressure and temperature. By making calculations using the experimental data we can identify the sample as having monoclinic structure at ambient conditions and a = 12.633 Å, b = 19.687 Å, c = 3.052 Å,  $\beta = 105.464^{\circ}$ . The unit-cell volume is 731.50 Å<sup>3</sup>. (The data were calculated using the 'Powder X' software, which uses 'Treor90' to index the powder x-ray diffraction patterns.)

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<sup>&</sup>lt;sup>3</sup> Author to whom any correspondence should be addressed.



Figure 1. The energy-dispersive x-ray diffraction pattern of Ga<sub>2</sub>O<sub>3</sub> under normal pressure.

A Mao–Bell-type diamond anvil cell (DAC) was used to generate high pressure. Using the cell, we carried out *in situ* high-pressure x-ray diffraction studies on Ga<sub>2</sub>O<sub>3</sub> using synchrotron radiation. The whole experiment took place in the Institute of High Energy Physics, Academia Sinica, in Beijing. The padding material was L605 (stainless steel). The padding material had been hardened by application of 10 GPa pressure. The radius of the sample container is about 200  $\mu$ m. The medium used to transfer pressure was a mixture of methyl alcohol, ethanol and water. Each spectrum was collected for 600–900 s. The position of the (111) diffraction peak of Pt and the formula *Ed* (keV nm) = 0.619 927/sin  $\theta$  was used to determine an internal pressure standard. The diffraction angle  $\theta = 8.244^{\circ}$ .

#### 3. Results and discussion

Figure 2 shows the spectra for energy-dispersive x-ray diffraction of  $Ga_2O_3$  at various pressures. One observes that all the diffraction peaks shift to higher energy with increasing pressure. At about 2.25 GPa, we identified the peak between Pt(111) and Pt(222) as being induced by padding material. So we relocated the cell and the peak disappeared.

At about 13.3 GPa, peaks 1 and 2 disappear. New peaks (3, 6 and 7) appear. As the pressure gets larger, these peaks become ever clearer. They are stable over the pressure range ( $\leq 38$  GPa) under study.

Figure 3 shows the *d*-spacing of  $Ga_2O_3$  under various pressures. It shows clearly the transformations of the peaks. After 13.3 GPa, the slopes of these lines has changed. Line 5 is broken; so we think that the new peaks 4 and 5 are no longer related to the old structure: they relate to the new phase that appears. Because the slope after 13.3 GPa is smaller, we can assume that the new structure is harder. Because  $\gamma$ -Ga<sub>2</sub>O<sub>3</sub> is definitely a harder material with cubic structure, we might think that Ga<sub>2</sub>O<sub>3</sub> became  $\gamma$ -Ga<sub>2</sub>O<sub>3</sub> after 13.3 GPa. But our calculations, with about 10% error, show that the new structure cannot be cubic. So the new structure is most probably tetragonal  $\beta$ -Ga<sub>2</sub>O<sub>3</sub>.

So we used the 'PowerX' software once again to process the data. We picked out the six strongest peaks at 37 GPa to index the powder x-ray diffraction patterns (see table 1).

The results from the software analysis show that the new structure is tetragonal with a = 8.097 Å, b = 8.097 Å and c = 10.062 Å. This result also explains why the new structure is very close to cubic structure.



Figure 2. Energy-dispersive x-ray diffraction patterns of Ga<sub>2</sub>O<sub>3</sub> under various pressures.



Figure 3. The *d*-spacing of Ga<sub>2</sub>O<sub>3</sub> under various pressures.

Table 1. The six strongest peaks at 37 GPa.

d-spacing (Å)	Intensity
3.41	174
2.58	645
2.46	1000
2.19	813
1.90	428
1.78	241

# 4. Conclusions

The experiments show that, in the energy-dispersive x-ray diffraction patterns of Ga<sub>2</sub>O<sub>3</sub>, with increasing pressure, the old peaks disappear (at about 13.3 GPa) and new peaks appear. A structural transition becomes apparent. Monoclinic  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub> turns into tetragonal  $\beta$ -Ga<sub>2</sub>O<sub>3</sub>. The new tetragonal structure ( $\beta$ -Ga<sub>2</sub>O<sub>3</sub>) can remain stable over the pressure range ( $\leq$ 38 GPa) of our research.

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